



Generalized quadrature formulae for analytic functions [☆]

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Dedicated to Prof. Dobrilo Đ. Tošić on the Occasion of his Eightieth Birth Anniversary.

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ABSTRACT

A kind of generalized quadrature formulae of maximal degree of precision for numerical integration of analytic functions is considered. Precisely, a general weighted quadrature of Birkhoff–Young type with $4n + 3$ nodes and degree of precision $6n + 5$ is studied. Its nodes are characterized by an orthogonality relation and a general numerical method for their computation is given. Special cases and numerical results are also included.

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1. Introduction

Recently Acharya et al. [1] have considered numerical approximation of integrals

$$I(w;f) = \int_L f(z)dz, \quad (1.1)$$

over a directed line segment L from the point $z_0 - h$ to the point $z_0 + h$ in the complex plane \mathbb{C} , where f is an analytic function in the disk

$$\Omega = \{z \in \mathbb{C} : |z - z_0| \leq r, r > |h|\}, \quad (1.2)$$

by means of a 7-point quadrature formula of the form

$$Q_7(w;f) = Af(z_0) + B[f(z_0 + th) + f(z_0 - th)] + C[f(z_0 + h) + f(z_0 - h)] + D[f(z_0 + ih) + f(z_0 - ih)],$$

where t is a positive parameter different from 1. Such a formula is exact for all odd degree monomials $f(z) = (z - z_0)^{2k+1}$. In order that the formula is also exact for even monomials $f(z) = (z - z_0)^{2k}$, $k = 0, 1, 2, 3$, the authors in [1] determine coefficients in $Q_7(w;f)$ as functions on t ,

$$(A, B, C, D) = \frac{h}{105} \left(\frac{8(21t^2 - 5)}{t^2}, \frac{20}{t^2(1 - t^4)}, \frac{2(9 - 14t^2)}{1 - t^2}, \frac{3 - 7t^2}{1 + t^2} \right),$$

and obtain a general formula of degree of precision at least seven for any finite positive value of the real parameter $t \neq 1$. Letting $t \rightarrow \infty$, this rule reduces to the well-known five-point Birkhoff–Young formula of fifth degree precision [3], for which

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$A = 8h/5$, $(B = 0)$, $C = 4h/15$, $D = -h/15$, and its remainder term $R_5^{BY}(w;f)$ can be estimated as (see [18] or Davis and Rabinowitz [5, p. 136])

$$|R_5^{BY}(w;f)| \leq \frac{|h|^7}{1890} \max_{z \in S} |f^{(6)}(z)|,$$

where S denotes the square with vertices $z_0 + i^k h$, $k = 0, 1, 2, 3$. The case of $Q_7(w;f)$ with $h = 1$ was considered in [17]. This kind of Birkhoff–Young quadrature formulae has been investigated by several authors [7,16,15,12–14]. Such quadrature formulae can also be used to integrate real harmonic functions (see [3]).

We mention here that Lyness and Delves [8] and Lyness and Moler [9], and later Lyness [10], developed formulae for numerical integration and numerical differentiation of complex functions.

By an analysis of the remainder term for the general 7-point quadrature formula $Q_7(w;f)$ with respect to the parameter t , it can be obtained a quadrature rule of the maximal precision nine for $t = \sqrt{7/15}$ (see [1]). Some other rules of degree of precision seven can be also derived.

However, with a little modification of $Q_7(w;f)$ we can obtain a modified 7-point quadrature rule $Q_7^M(w;f)$ of degree precision eleven. Furthermore, using such an approach we derive a general $(4n + 3)$ -point quadrature formula of the maximal degree of precision for a weighted integral.

The paper is organized as follows. In Section 2 we give the modified quadrature formula $Q_7^M(w;f)$. Section 3 is devoted to a general weighted quadrature of Birkhoff–Young type with $4n + 3$ nodes and degree of precision $6n + 5$. The nodes of such quadratures are characterized by an orthogonality relation. The corresponding weight coefficients of quadratures are given in Section 4. A general numerical method for determining nodes of such quadratures of maximal degree of precision is discussed in Section 5, including numerical results.

2. The modified quadrature formula $Q_7^M(w;f)$

For numerical calculating of the integral (1.1) of an analytic function in the disk (1.2), in this section we consider a modification of the quadrature formula $Q_7(w;f)$ in the following form

$$Q_7^M(w;f) = Af(z_0) + B[f(z_0 + th) + f(z_0 - th)] + C[f(z_0 + \ell h) + f(z_0 - \ell h)] + D[f(z_0 + i\ell h) + f(z_0 - i\ell h)],$$

where t and ℓ are mutually different positive parameters. In this case, from the corresponding system of equations

$$\begin{aligned} \frac{1}{2}A + B + C + D &= h, \\ Bt^2 + C\ell^2 - D\ell^2 &= \frac{h}{3}, \\ Bt^4 + C\ell^4 + D\ell^4 &= \frac{h}{5}, \\ Bt^6 + C\ell^6 - D\ell^6 &= \frac{h}{7}, \end{aligned}$$

we get

$$\begin{aligned} A &= \frac{2h}{105} \cdot \frac{105t^2\ell^4 - 21t^2 - 35\ell^4 + 15}{t^2\ell^4}, & B &= \frac{h}{21} \cdot \frac{7\ell^4 - 3}{t^2(\ell^4 - t^4)}, \\ C &= \frac{h}{210} \cdot \frac{35t^2\ell^2 + 21t^2 - 21\ell^2 - 15}{\ell^4(t^2 - \ell^2)}, & D &= \frac{h}{210} \cdot \frac{-35t^2\ell^2 + 21t^2 + 21\ell^2 - 15}{\ell^4(\ell^2 + t^2)}, \end{aligned}$$

where $0 < t$, $\ell < 1$, $t \neq \ell$.

It is pertinent to note that the modified quadrature formula $Q_7^M(w;f)$ boils down to the seventh degree rule due to Acharya et al. [1] for $\ell = 1$, and the modified Birkhoff–Young rule due to Tošić [16] for $\ell = (3/7)^{1/4}$ and $t \rightarrow \infty$.

Now, the error-term $R_7^M(w;f) = I(w;f) - Q_7^M(w;f)$ for $f(z) = (z - z_0)^{2k}$, $k = 4, 5, 6$, reduces to

$$\begin{aligned} R_7((z - z_0)^8) &= \frac{2h^9}{315} [21\ell^4(5t^2 - 3) - 45t^2 + 35], \\ R_7((z - z_0)^{10}) &= \frac{2h^{11}}{231} [11\ell^4(7t^4 - 3) - 33t^4 + 21], \\ R_7((z - z_0)^{12}) &= \frac{2h^{13}}{1365} [91\ell^8(5t^2 - 3) + 65\ell^4t^2(7t^4 - 3) - 195t^6 + 105]. \end{aligned}$$

Finally, from

$$R_7((z - z_0)^8) = 0, \quad R_7((z - z_0)^{10}) = 0,$$

we obtain two solutions for parameters:

$$\ell^A = \frac{5}{693}(57 + 4\sqrt{102}), \quad t^2 = \frac{1}{77}(45 - 2\sqrt{102}), \quad (2.1)$$

and

$$\ell^A = \frac{5}{693}(57 - 4\sqrt{102}), \quad t^2 = \frac{1}{77}(45 + 2\sqrt{102}). \quad (2.2)$$

Thus, there exist two modified 7-point quadratures $Q_{7,\nu}^M(w;f)$, $\nu = 0, 1$, of the maximal degree of precision eleven, with the following parameters:

$$\begin{aligned} Q_{7,0}^M(w;f) : \quad \ell &= \sqrt[4]{\frac{5}{693}(57 + 4\sqrt{102})} \approx 0.9155808999196944, \\ t &= \sqrt{\frac{1}{77}(45 - 2\sqrt{102})} \approx 0.5675304228160498, \\ A &= \frac{256(198 - \sqrt{102})}{77175} = 0.6232915676809758h, \\ B &= \frac{2939400 + 116087\sqrt{102}}{8680644} = 0.4736769794706059h, \\ C &= 0.2151573287932331h, \quad D = -0.0004800921043269324h, \end{aligned}$$

and

$$\begin{aligned} Q_{7,1}^M(w;f) : \quad \ell &= \sqrt[4]{\frac{5}{693}(57 - 4\sqrt{102})} \approx 0.5883004297385740, \\ t &= \sqrt{\frac{1}{77}(45 + 2\sqrt{102})} \approx 0.9201849748878780, \\ A &= \frac{256(198 + \sqrt{102})}{77175} = 0.6902944381499280h, \\ B &= \frac{2939400 - 116087\sqrt{102}}{8680644} = 0.2035538803596094h, \\ C &= 0.4582083249363621h, \quad D = -0.006909424370935494h. \end{aligned}$$

Expanding an analytic function in Taylor series

$$f(z) = \sum_{k=0}^{+\infty} \frac{f^{(k)}(z_0)}{k!} (z - z_0)^k, \quad z \in \Omega,$$

and having in mind that $Q_7^M(w;f)$ has the maximal degree of precision eleven, the remainder is

$$R_7^M(w;f) = \frac{f^{(12)}(z_0)}{12!} R_7^M((z - z_0)^{12}) + \dots$$

For the obtained parameters (2.1) and (2.2), its dominant error term reduces to

$$R_{7,0}^M(w;f) \approx \frac{256(516 + 13\sqrt{102})}{3361743} \cdot \frac{h^{13}}{13!} f^{(12)}(z_0) \approx 7.92 \times 10^{-12} h^{13} f^{(12)}(z_0),$$

and

$$R_{7,1}^M(w;f) \approx \frac{256(516 - 13\sqrt{102})}{3361743} \cdot \frac{h^{13}}{13!} f^{(12)}(z_0) \approx 4.70 \times 10^{-12} h^{13} f^{(12)}(z_0),$$

respectively. As we can see the second formula is slightly more accurate than the first one.

3. The generalized weighted quadrature formula $Q_{4n+3}^M(w;f)$

Let $w : (-1, 1) \rightarrow \mathbb{R}^+$ be an even positive weight function, for which all moments $\mu_k = \int_{-1}^1 z^k w(z) dz$, $k = 0, 1, \dots$, exist. Without loss of generality, in this section we consider a weighted integration over $L = [-1, 1]$ for analytic functions in the unit disk $\Omega = \{z : |z| \leq 1\}$ by

$$I(w;f) = \int_{-1}^1 f(z)w(z)dz = Q_{4n+3}(w;f) + R_{4n+3}(w;f), \quad (3.1)$$

where $Q_{4n+3}(w; f)$ is the $(4n+3)$ -point quadrature formula of interpolatory type with nodes at the zeros of a monic polynomial of degree $4n+3$ ($n \in \mathbb{N}$),

$$\Omega_{4n+3}(z) = z(z^2 - r_0)p_n(z^4) = z(z^2 - r_0) \prod_{k=1}^n (z^4 - r_k), \quad (3.2)$$

where $0 < r_1 < \dots < r_n < 1$, $r_0 \in (0, 1)$, and $r_0 \neq r_k$ for each $k = 1, \dots, n$. Here, $R_{4n+3}(w; f)$ is the corresponding remainder term.

According to (3.2) the quadrature formula in (3.1) has the form

$$Q_{4n+3}(w; f) = Af(0) + B[f(x_0) + f(-x_0)] + \sum_{k=1}^n \{C_k[f(x_k) + f(-x_k)] + D_k[f(ix_k) + f(-ix_k)]\}, \quad (3.3)$$

where $x_0 = \sqrt{r_0}$ and $x_k = \sqrt[4]{r_k}$, $k = 1, \dots, n$.

Theorem 3.1. For any $n \in \mathbb{N}$ there exist interpolatory quadratures $Q_{4n+3}(w; f)$ with a maximal degree of precision $d = 6n + 5$. The nodes of such quadratures $Q_{4n+3}^M(w; f)$ are characterized by the following orthogonality relation

$$\int_{-1}^1 \pi_{2k}(z) z^2 (z^2 - r_0) p_n(z^4) w(z) dz = 0, \quad k = 0, 1, \dots, n, \quad (3.4)$$

where $\{\pi_\nu\}_{\nu \in \mathbb{N}_0}$ is a sequence of orthogonal polynomials with respect to the weight function on $(-1, 1)$.

Proof. Let \mathcal{P}_d denotes the set of algebraic polynomials of degree at most d . For a given $n \in \mathbb{N}$, suppose that $f \in \mathcal{P}_d$, where $d \geq 4n + 3$. Then, it can be expressed in the form

$$f(z) = u(z)\Omega_{4n+3}(z) + v(z) = u(z)z(z^2 - r_0)p_n(z^4) + v(z),$$

where $u \in \mathcal{P}_{d-4n-3}$ and $v \in \mathcal{P}_{4n+2}$. Applying (3.1), we get

$$I(w; f) = \int_{-1}^1 u(z)z(z^2 - r_0)p_n(z^4)w(z)dz + I(v).$$

Since this quadrature is of interpolatory type we have that $I(w; v) = Q_{4n+3}(v)$ and also $v(z) = f(z)$ at the zeros of the polynomial Ω_{4n+3} . Therefore, $Q_{4n+3}(w; v) = Q_{4n+3}(w; f)$, so that for each $f \in \mathcal{P}_d$ we have

$$I(w; f) = \int_{-1}^1 u(z)z(z^2 - r_0)p_n(z^4)w(z)dz + Q_{4n+3}(w; f).$$

It is clear that the quadrature formula $Q_{4n+3}(w; f)$ becomes $Q_{4n+3}^M(w; f)$, i.e., it has a maximal degree of precision, if and only if

$$\int_{-1}^1 u(z)z(z^2 - r_0)p_n(z^4)w(z)dz = 0 \quad (3.5)$$

for a maximal degree of polynomials $u \in \mathcal{P}_{d-4n-3}$. Evidently, (3.5) is true for every even polynomial. Taking u as an odd polynomial $u(z) = zh(z^2)$, where $h \in \mathcal{P}_n$, the previous “orthogonality conditions” can be represented as

$$\int_{-1}^1 h(z^2)z^2(z^2 - r_0)p_n(z^4)w(z)dz = 0, \quad h \in \mathcal{P}_n. \quad (3.6)$$

Since the maximal degree of the polynomial $u \in \mathcal{P}_{d-4n-3}$ is

$$d_{\max} - 4n - 3 = 1 + 2n + 1,$$

we conclude that the maximal degree of precision of such a quadrature $Q_{4n+3}(w; f)$ is $d_{\max} = 6n + 5$, i.e., $Q_{4n+3}(w; f) = Q_{4n+3}^M(w; f)$.

Introducing the inner product in a usual way as

$$(f, g) = \int_{-1}^1 f(z)g(z)w(z)dz,$$

the last orthogonality conditions (3.6) can be expressed in terms of orthogonal polynomials $\{\pi_\nu\}_{\nu \in \mathbb{N}_0}$ with respect to this inner product in the form $(z\pi_{2k}, \Omega_{4n+3}) = 0$, $0 \leq k \leq n$, i.e., (3.4). \square

According to (3.2), the polynomial $(z^4 - rz^2)p_n(z^4)$ can be expressed in the form

$$(z^4 - rz^2)p_n(z^4) = \sum_{j=0}^n (-1)^j \sigma_j (z^{4(n-j+1)} - rz^{2(2n-2j+1)}), \quad (3.7)$$

where σ_j are the so-called *elementary symmetric functions*, defined by

$$\sigma_j = \sum_{(k_1, \dots, k_j)} r_{k_1} \cdots r_{k_j}, \quad j = 1, \dots, n,$$

and the summation is performed over all combinations (k_1, \dots, k_j) of the basic set $\{1, \dots, n\}$. Thus,

$$\sigma_1 = r_1 + r_2 + \cdots + r_n, \quad \sigma_2 = r_1 r_2 + \cdots + r_{n-1} r_n, \dots, \quad \sigma_n = r_1 r_2 \cdots r_n,$$

and for the convenience we put $\sigma_0 = 1$. Also, we put r instead of r_0 .

Using the orthogonality conditions (3.4) and the expansion (3.7) we get the following system of nonlinear equations

$$f_k \equiv \sum_{j=0}^n (-1)^j \sigma_j \{s_{k,2n-2j+2} - r s_{k,2n-2j+1}\} = 0, \quad k = 0, 1, \dots, n, \quad (3.8)$$

with respect to unknowns $r, \sigma_1, \dots, \sigma_n$, or equivalently to $r_k, k = 0, 1, \dots, n$, where $r = r_0$ and $s_{k,j} = (\pi_{2k}, z^{2j}), k, j \geq 0$.

Introducing the notations $\boldsymbol{\sigma} = [\sigma_1 \ \sigma_2 \ \cdots \ \sigma_n]^T$,

$$\mathbf{A} = \mathbf{A}(r) = [a_{k,j}]_{k=0, j=1}^{n,n}, \quad \mathbf{b} = \mathbf{b}(r) = [b_0 \ b_1 \ \cdots \ b_n]^T,$$

where

$$a_{k,j} = (-1)^{j-1} [s_{k,2n-2j+2} - r s_{k,2n-2j+1}], \quad b_k = s_{k,2n+2} - r s_{k,2n+1}, \quad (3.9)$$

the system of $n + 1$ nonlinear equations (3.8) can be written in the matrix form

$$\mathbf{A}\boldsymbol{\sigma} = \mathbf{b}. \quad (3.10)$$

We have seen that the problem for $n = 1$ (and $w(z) = 1$) has two solutions. Numerical experiments show that for an arbitrary n , the number of solutions is $n + 1$. This hypothesis can be checked numerically for some reasonable values of n (e.g. $n \leq 10$) in the following way.

If we take a fixed value of $r \in (0, 1)$, then the overdetermined system of $n + 1$ linear equations

$$\mathbf{A}(r)\boldsymbol{\sigma} = \mathbf{b}(r), \quad (3.11)$$

with n unknowns $\boldsymbol{\sigma} = (\sigma_1, \sigma_2, \dots, \sigma_n)$, can be solved as a least squares problem (in the 2-norm)

$$\min_{\boldsymbol{\sigma}} \|\mathbf{A}(r)\boldsymbol{\sigma} - \mathbf{b}(r)\|_2 = \|\mathbf{A}(r)\hat{\boldsymbol{\sigma}} - \mathbf{b}(r)\|_2,$$

where the vector $\mathbf{A}(r)\hat{\boldsymbol{\sigma}} - \mathbf{b}(r)$ is the corresponding *least squares residual* and the solution $\hat{\boldsymbol{\sigma}}$ can be expressed in terms of Moore–Penrose inverse.

Only when for some $r = \hat{r}_0$, the vector $\mathbf{A}(r)\hat{\boldsymbol{\sigma}} - \mathbf{b}(r)$ becomes zero, we can identify the existence of a solution $(\hat{r}_0, \hat{\boldsymbol{\sigma}}) = (\hat{r}_0, \hat{\sigma}_1, \dots, \hat{\sigma}_n)$ of our original (nonlinear) system of Eqs. (3.10).

Consider now the case $w(z) = 1$. Following [2, §10.10] we get

$$(P_{2k}, z^{2j}) = k! \binom{j}{k} \frac{\Gamma(j+1/2)}{\Gamma(k+j+3/2)} = k! \binom{j}{k} \prod_{v=0}^k \frac{2}{2j+2v+1}, \quad (3.12)$$

where $P_{2k}(z)$ is the Legendre polynomial of degree $2k$. Taking (3.12) instead of $s_{k,j}$, the corresponding norm $\|\mathbf{A}(r)\hat{\boldsymbol{\sigma}} - \mathbf{b}(r)\|_2$ as a function of r is presented in Fig. 1 for $1 \leq n \leq 4$ and $n = 10$.

As we can see, for the weight function $w(z) = 1$ and these values of n , the 2-norm vanishes only at $n + 1$ points in $(0, 1)$, which means that for a given n , there exist $n + 1$ different quadratures $Q_{4n+3, \nu}^M, \nu = 0, 1, \dots, n$, each of degree of precision $d_{\max} = 6n + 5$.

In a general case, if we have a solution of our nonlinear problem, say \hat{r}_0 , then in order to construct the corresponding quadrature formula $Q_{4n+3}^M(w; f)$ for such a $r = \hat{r}_0$, we should solve a system of n linear equations from (3.11) in order to get the values $(\hat{\sigma}_1, \dots, \hat{\sigma}_n)$, and then the zeros $(\hat{r}_1, \dots, \hat{r}_n)$ by solving the equation

$$(z - \hat{r}_1) \cdots (z - \hat{r}_n) = z^n - \hat{\sigma}_1 z^{n-1} + \hat{\sigma}_2 z^{n-2} - \cdots + (-1)^n \hat{\sigma}_n = 0.$$

Then, the nodes in the corresponding quadrature formula (3.3) are

$$x_0 = \sqrt{\hat{r}_0} \quad \text{and} \quad x_k = \sqrt[4]{\hat{r}_k}, \quad k = 1, \dots, n.$$

A determination of the weight coefficients A, B, C_k and $D_k, k = 1, \dots, n$, in the corresponding quadrature formula (3.3) is a linear problem and it is considered in the next section.

A general numerical method for solving our nonlinear problem (3.8) and finding all solutions for r_0 , for a given n and for an arbitrary weight function $w(z)$, is given in Section 5. Also, a general numerical method for calculating the necessary inner products $s_{k,j} = (\pi_{2k}, z^{2j}), 0 \leq k \leq j$, is given. Furthermore, analytic expressions for the generalized Gegenbauer weight are derived.

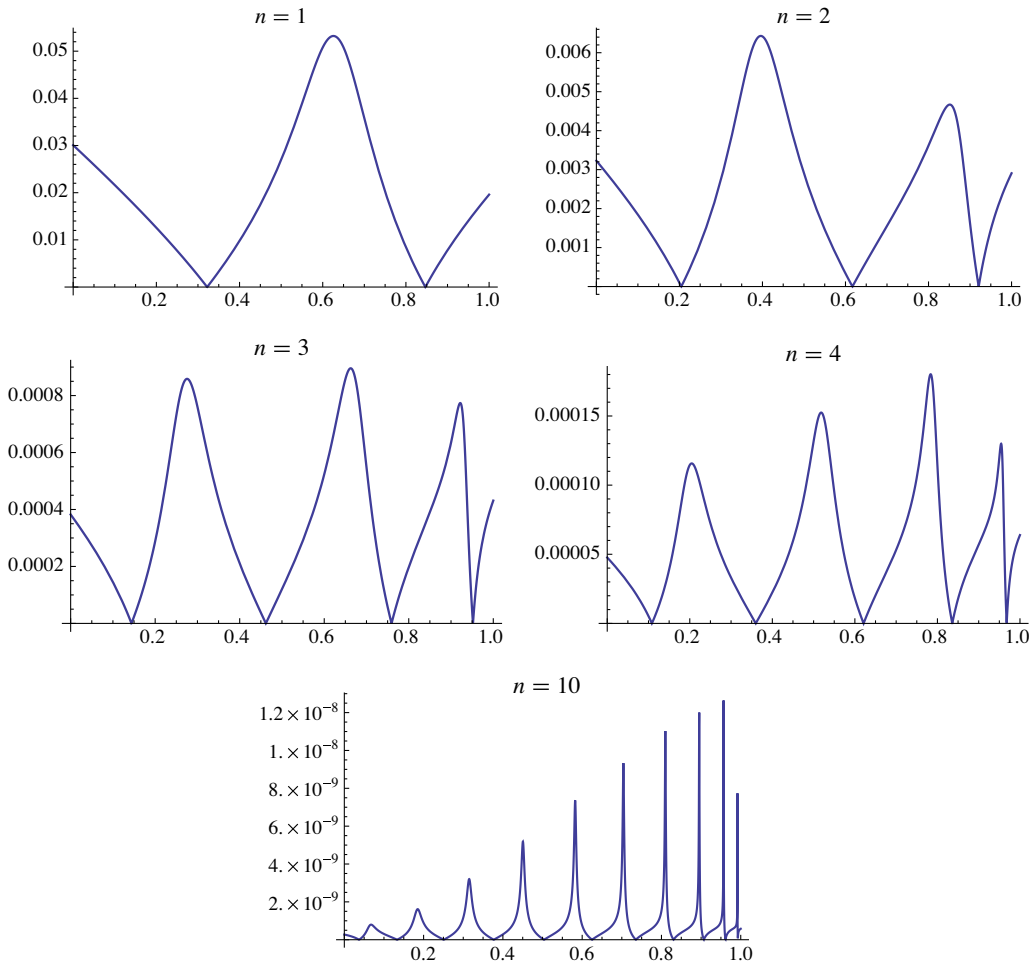


Fig. 1. The norm $\|A(r)\hat{\sigma} - b(r)\|_2$ as a function of r for $1 \leq n \leq 4$ and $n = 10$.

4. The weight coefficients in quadrature formula $Q_{4n+3}^M(w; f)$

In order to determine the weight coefficients W_ν in a quadrature formula of interpolatory type

$$Q_{4n+3}(w; f) = \sum_{z_\nu \in Z} W_\nu f(z_\nu),$$

we use the Lagrange polynomial constructed at the set of simple nodes $Z = \{z_\nu\}$. In our case,

$$Z = \{0, \pm x_0, \pm x_k, \pm ix_k, k = 1, \dots, n\},$$

where $x_0 = \sqrt{r_0}$ and $x_k = \sqrt[4]{r_k}$, $k = 1, \dots, n$, and the node polynomial is

$$\Omega_{4n+3}(z) = z(z^2 - r_0) \prod_{k=1}^n (z^4 - r_k).$$

The corresponding Lagrange polynomial is

$$L_{4n+3}(f; z) = \sum_{z_\nu \in Z} \frac{\Omega_{4n+3}(z)}{(z - z_\nu)\Omega'_{4n+3}(z_\nu)} f(z_\nu),$$

so that

$$W_\nu = \frac{1}{\Omega'_{4n+3}(z_\nu)} \int_{-1}^1 \frac{\Omega_{4n+3}(z)w(z)}{z - z_\nu} dz. \tag{4.1}$$

Theorem 4.1. Let r_0 and r_k , $k = 1, \dots, n$, be determined according to Theorem 3.1. Then the weight coefficients in the quadrature formula $Q_{4n+3}^M(w; f)$ with the maximal degree of precision $d = 6n + 5$ are given by

$$\begin{aligned} A &= \frac{-1}{r_0 p_n(0)} \int_{-1}^1 (z^2 - r_0) p_n(z^4) w(z) dz, \\ B &= \frac{1}{2r_0 p_n(r_0^2)} \int_{-1}^1 z^2 p_n(z^4) w(z) dz, \\ C_k &= \frac{1}{4r_k(\sqrt{r_k} - r_0) p_n'(r_k)} \int_{-1}^1 \frac{z^2(z^2 - r_0) p_n(z^4)}{z^2 - \sqrt{r_k}} w(z) dz, \quad k = 1, \dots, n, \\ D_k &= \frac{-1}{4r_k(\sqrt{r_k} + r_0) p_n'(r_k)} \int_{-1}^1 \frac{z^2(z^2 - r_0) p_n(z^4)}{z^2 + \sqrt{r_k}} w(z) dz, \quad k = 1, \dots, n, \end{aligned}$$

where $p_n(z) = \prod_{v=1}^n (z - r_v)$.

Proof. Let $x_0 = \sqrt{r_0}$ and $x_k = \sqrt[4]{r_k}$, $k = 1, \dots, n$. According to

$$\Omega'_{4n+3}(z) = (3z^2 - r_0) \prod_{v=1}^n (z^4 - r_v) + 4z^4(z^2 - r_0) \sum_{j=1}^n \prod_{v \neq j} (z^4 - r_v),$$

we have

$$\begin{aligned} \Omega'_{4n+3}(0) &= -r_0 \prod_{v=1}^n (-r_v) = -r_0 p_n(0), \\ \Omega'_{4n+3}(\pm x_0) &= 2r_0 \prod_{v=1}^n (r_0^2 - r_v) = 2r_0 p_n(r_0^2), \\ \Omega'_{4n+3}(\pm x_k) &= 4r_k(\sqrt{r_k} + r_0) \prod_{v \neq k} (r_k - r_v) = 4r_k(\sqrt{r_k} - r_0) p_n'(r_k), \\ \Omega'_{4n+3}(\pm i x_k) &= 4r_k(-\sqrt{r_k} - r_0) \prod_{v \neq k} (r_k - r_v) = -4r_k(\sqrt{r_k} + r_0) p_n'(r_k), \end{aligned}$$

where $k = 1, \dots, n$. Now, applying (4.1) and using notations for coefficients as in (3.3), we get desired results. \square

For analytic functions we can give an explicit formula for the interpolation error $E_{4n+3}(z) = f(z) - L_{4n+3}(f; z)$ in the form (see [11, pp. 55–56])

$$E_{4n+3}(z) = \frac{1}{2\pi i} \oint_{\Gamma} \frac{\Omega_{4n+3}(z)}{\Omega_{4n+3}(\zeta)} \frac{f(\zeta)}{\zeta - z} d\zeta \quad (z \in \text{int}\Gamma),$$

where Γ is a simple closed contour in \mathbb{C} , such that all interpolation nodes belong to $\text{int}\Gamma$. Then, the remainder term in (3.1) can be expressed in the integral form

$$R_{4n+3}(w; f) = \frac{1}{2\pi i} \oint_{\Gamma} \frac{f(\zeta)}{\Omega_{4n+3}(\zeta)} \left(\int_{-1}^1 \frac{\Omega_{4n+3}(z) w(z)}{\zeta - z} dz \right) d\zeta. \quad (4.2)$$

Some estimate of (4.2) will be given elsewhere.

5. Numerical methods for constructing quadratures

In the sequel we need the inner products $s_{k,j} = (\pi_{2k}, z^{2j})$, $0 \leq k \leq j$. For $k = 0$ these products reduce to the moments of the weight function w , i.e.,

$$s_{0,j} = (1, z^{2j}) = \int_{-1}^1 z^{2j} w(z) dz = \mu_{2j}, \quad j = 0, 1, \dots \quad (5.1)$$

First, we give an analytic expression of $s_{k,j}$ for a wide class of weight functions, and after that we introduce a general numerical method for easy calculation of $s_{k,j}$ for every even weight function.

5.1. Analytic expression of $s_{k,j}$ for the generalized Gegenbauer weight

We consider the so-called generalized Gegenbauer weight function defined by $w(z) = |z|^\gamma (1 - z^2)^\alpha$, $\gamma, \alpha > -1$, on $(-1, 1)$. The monic polynomials $W_v^{(\alpha, \beta)}(z)$, $v = 0, 1, \dots$, orthogonal with respect to this weight function, where $\beta = (\gamma - 1)/2$, were introduced by Lašćenov [6] (see, also, Chihara [4, pp. 155–156] and Mastroianni and Milovanović [11, pp. 147–148]). These

polynomials can be expressed in terms of the Jacobi polynomials $P_v^{(\alpha, \beta)}(z)$, $v = 0, 1, \dots$, which are orthogonal on $(-1, 1)$ with respect to the weight function $w^{(\alpha, \beta)}(z) = (1-z)^\alpha(1+z)^\beta$, $\alpha, \beta > -1$. Namely,

$$\begin{aligned} W_{2k}^{(\alpha, \beta)}(z) &= \frac{k!}{(k + \alpha + \beta + 1)_k} P_k^{(\alpha, \beta)}(2z^2 - 1), \\ W_{2k+1}^{(\alpha, \beta)}(z) &= \frac{k!}{(k + \alpha + \beta + 2)_k} z P_k^{(\alpha, \beta+1)}(2z^2 - 1). \end{aligned} \quad (5.2)$$

Notice that $W_{2k+1}^{(\alpha, \beta)}(z) = zW_{2k}^{(\alpha, \beta+1)}(z)$. These polynomials satisfy the following three-term recurrence relation

$$\begin{aligned} W_{v+1}^{(\alpha, \beta)}(z) &= zW_v^{(\alpha, \beta)}(z) - \beta_v W_{v-1}^{(\alpha, \beta)}(z), \quad v = 0, 1, \dots, \\ W_{-1}^{(\alpha, \beta)}(z) &= 0, \quad W_0^{(\alpha, \beta)}(z) = 1, \end{aligned}$$

where

$$\beta_{2k} = \frac{k(k + \alpha)}{(2k + \alpha + \beta)(2k + \alpha + \beta + 1)}, \quad \beta_{2k-1} = \frac{(k + \beta)(k + \alpha + \beta)}{(2k + \alpha + \beta - 1)(2k + \alpha + \beta)},$$

for $k = 1, 2, \dots$, except when $\alpha + \beta = -1$; then $\beta_1 = (\beta + 1)/(\alpha + \beta + 2)$.

Now, we want to find an explicit expression for the products

$$s_{k,j} = (W_{2k}^{(\alpha, \beta)}, z^{2j}) = \int_{-1}^1 W_{2k}^{(\alpha, \beta)}(z) z^{2j} |z|^{2\beta+1} (1-z^2)^\alpha dz, \quad 0 \leq k \leq j. \quad (5.3)$$

Lemma 5.1. Let $\alpha, \beta > -1$. Then the products defined in (5.3) are

$$s_{k,j} = \frac{k!}{(k + \alpha + \beta + 1)_k} \binom{j}{k} \frac{\Gamma(k + \alpha + 1)\Gamma(j + \beta + 1)}{\Gamma(k + j + \alpha + \beta + 2)} \quad (5.4)$$

for $0 \leq k \leq j$. Otherwise, $s_{k,j} = 0$.

In order to prove this lemma we need an auxiliary result.

Lemma 5.2. For $\alpha, \beta > -1$ and $j \in \mathbb{N}_0$ the expansion in Jacobi polynomials $P_v^{(\alpha, \beta)}(t)$, $v = 0, 1, \dots, j$,

$$(1+t)^j = 2^j \Gamma(j + \beta + 1) \sum_{v=0}^j v! \binom{j}{k} \frac{(2v + \alpha + \beta + 1)\Gamma(v + \alpha + \beta + 1)}{\Gamma(v + \beta + 1)\Gamma(v + j + \alpha + \beta + 2)} P_v^{(\alpha, \beta)}(t)$$

holds.

The proof of this expansion can be given by induction in j , having in mind that

$$\begin{aligned} (1+t)P_v^{(\alpha, \beta)}(t) &= \frac{2(v+1)(v + \alpha + \beta + 1)}{(2v + \alpha + \beta + 1)(2v + \alpha + \beta + 2)} P_{v+1}^{(\alpha, \beta)}(t) + \left(1 + \frac{\beta^2 - \alpha^2}{(2v + \alpha + \beta)(2v + \alpha + \beta + 1)}\right) P_v^{(\alpha, \beta)}(t) \\ &\quad + \frac{2(v + \alpha)(v + \beta)}{(2v + \alpha + \beta)(2v + \alpha + \beta + 1)} P_{v-1}^{(\alpha, \beta)}(t). \end{aligned}$$

We mention that the corresponding expansion in Bateman and Erdélyi [2, pp. 212] has a mistake.

Proof of Lemma 5.2. According to (5.3) and (5.2) we have

$$s_{k,j} = \frac{2k!}{(k + \alpha + \beta + 1)_k} \int_0^1 P_k^{(\alpha, \beta)}(2z^2 - 1) z^{2j+2\beta+1} (1-z^2)^\alpha dz.$$

Changing the variables $t = 2z^2 - 1$ it reduces to

$$s_{k,j} = \frac{k!}{2^{\alpha+\beta+j+1} (k + \alpha + \beta + 1)_k} \int_{-1}^1 P_k^{(\alpha, \beta)}(t) (1-t)^\alpha (1+t)^{\beta+j} dt.$$

Now, using the expansion from Lemma 5.2 and the orthogonality of Jacobi polynomials,

$$(P_k^{(\alpha, \beta)}, P_v^{(\alpha, \beta)}) = \|P_k^{(\alpha, \beta)}\|^2 \delta_{k,v} = \frac{2^{\alpha+\beta+1} \Gamma(k + \alpha + 1)\Gamma(k + \beta + 1)}{k!(2k + \alpha + \beta + 1)\Gamma(k + \alpha + \beta + 1)} \delta_{k,v},$$

where $\delta_{k,v}$ is Kronecker's delta, we obtain

$$s_{k,j} = \frac{k!2^j\Gamma(j+\beta+1)k!\binom{j}{k}}{2^{\alpha+\beta+j+1}(k+\alpha+\beta+1)_k} \frac{(2k+\alpha+\beta+1)\Gamma(k+\alpha+\beta+1)}{\Gamma(k+\beta+1)\Gamma(k+j+\alpha+\beta+2)} \|P_k^{(\alpha,\beta)}\|^2,$$

i.e., (5.4). For $k > j$, because of orthogonality, $s_{k,j} = 0$ \square

We mention some special cases:

1° $w(z) = 1$, i.e., $\alpha = 0$, $\gamma = 0$ ($\beta = -1/2$): In this Legendre case we have

$$s_{k,j} = \frac{(k!)^2}{(k+1/2)_k} \binom{j}{k} \frac{\Gamma(j+1/2)}{\Gamma(k+j+3/2)}.$$

Compared with (3.12), the additional factor $k!/(k+1/2)_k$ ($= 1/a_{2k}$) comes from the leading coefficient in the Legendre polynomial $P_{2k}(z) = a_{2k}z^{2k} + \text{terms of lower degree}$.

2° $w(z) = 1/\sqrt{1-z^2}$, i.e., $\alpha = -1/2$, $\gamma = 0$ ($\beta = -1/2$): In the Chebyshev case of the first kind, the inner product (5.4) reduces to

$$s_{k,j} = \frac{\pi}{2^{2j+2k-1}} \binom{2j}{j-k}.$$

3° $w(z) = \sqrt{1-z^2}$, i.e., $\alpha = 1/2$, $\gamma = 0$ ($\beta = -1/2$): In the Chebyshev case of the second kind, it becomes

$$s_{k,j} = \frac{\pi}{2^{2j+2k+1}} \frac{2k+1}{2j+1} \binom{2j+1}{j-k}.$$

4° $w(z) = (1-z^2)^\alpha$, i.e., $\alpha > -1$, $\gamma = 0$ ($\beta = -1/2$): In this Gegenbauer case, we have

$$s_{k,j} = \frac{k!}{(k+\alpha+1)_k} \binom{j}{k} \frac{\Gamma(k+\alpha+1)\Gamma(j+1/2)}{\Gamma(k+j+\alpha+3/2)}.$$

5° $w(z) = |z|$, i.e., $\alpha = 0$, $\gamma = 1$ ($\beta = 0$): In this case

$$s_{k,j} = \frac{1}{k+j+1} \frac{\binom{2j}{j-k}}{\binom{2k}{k}\binom{2j}{j}}.$$

5.2. General numerical method for calculating $s_{k,j}$

It is well-known that monic polynomials $\{\pi_v\}_{v \in \mathbb{N}_0}$ orthogonal with respect to an even weight function satisfy the three-term recurrence relation of the form

$$\pi_{v+1}(z) = z\pi_v(z) - \beta_v\pi_{v-1}(z), \quad v = 0, 1, \dots, \quad (5.5)$$

with $\pi_0(z) = 1$ and $\pi_{-1}(z) = 0$. It is convenient to put $\beta_0 = \mu_0$.

Lemma 5.3. Let β_v , $v \geq 0$, be recursion coefficients in (5.5) for polynomials orthogonal with respect to the even weight function w on $(-1, 1)$, with the moments $\mu_v = \int_{-1}^1 z^v w(z) dz$, $v \geq 0$. For the inner products $s_{k,j} = (\pi_{2k}, z^{2j})$ the following recurrence relation

$$s_{k,j+1} = s_{k+1,j} + (\beta_{2k} + \beta_{2k+1})s_{k,j} + \beta_{2k}\beta_{2k-1}s_{k-1,j} \quad (5.6)$$

holds, with $s_{0,j} = \mu_{2j}$, $j = 0, 1, \dots$, and $s_{k,j} = 0$ for $k > j$.

Remark 5.4. Coefficients from the relation (5.6) appear in the recurrence relation for polynomials $\{\pi_{2k}(\sqrt{t})\}_{k \in \mathbb{N}_0}$ orthogonal with respect to the weight function $w(\sqrt{t})/\sqrt{t}$ on $(0, 1)$ (see [11, pp. 101–103]).

Proof of Lemma 5.3. Because of orthogonality, it is clear that $s_{k,j} = (\pi_{2k}, z^{2j}) = 0$ for $k > j$. When $k = 0$, for the boundary values $s_{0,j}$ we have (5.1). For diagonal elements we have

$$s_{j,j} = (\pi_{2j}, z^{2j}) = \|\pi_{2j}\|^2 = \prod_{v=0}^{2j} \beta_v.$$

In order to get (5.6), we start with the recurrence relation (5.5). Thus,

$$\begin{aligned} s_{k+1,j} &= \int_{-1}^1 \pi_{2k+2}(z)w(z)dz = \int_{-1}^1 [z\pi_{2k+1}(z) - \beta_{2k+1}\pi_{2k}(z)]z^{2j}w(z)dz \\ &= \int_{-1}^1 \pi_{2k+1}(z)z^{2j+1}w(z)dz - \beta_{2k+1} \int_{-1}^1 \pi_{2k}(z)z^{2j}w(z)dz = \int_{-1}^1 [z\pi_{2k}(z) - \beta_{2k}\pi_{2k-1}(z)]z^{2j+1}w(z)dz - \beta_{2k+1}s_{k,j} \\ &= s_{k,j+1} - \beta_{2k} \int_{-1}^1 z\pi_{2k-1}(z)z^{2j}w(z)dz - \beta_{2k+1}s_{k,j}. \end{aligned}$$

Expanding $z\pi_{2k-1}(z)$ as a linear combination of $\pi_{2k}(z)$ and $\pi_{2k-2}(z)$, we get

$$s_{k,j+1} = s_{k+1,j} + \beta_{2k} \int_{-1}^1 [\pi_{2k}(z) + \beta_{2k-1}\pi_{2k-2}(z)]z^{2j}w(z)dz + \beta_{2k+1}s_{k,j} = s_{k+1,j} + \beta_{2k}s_{k,j} + \beta_{2k}\beta_{2k-1}s_{k-1,j} + \beta_{2k+1}s_{k,j},$$

i.e., (5.6). \square

Fig. 2 displays the triangular array of the inner products $s_{k,j} = (\pi_{2k}, z^{2j})$, $0 \leq k \leq j$, and the computing stencil showing that the circled entry is computed in terms of the three other entries. The entries in the boxes are the known boundary values

$$s_{0,j} = \mu_{2j}, \quad s_{j,j} = \beta_0\beta_1 \cdots \beta_{2j}, \quad s_{j+1,j} = 0, \quad j = 0, 1, \dots, 2n + 2.$$

Zero entries $s_{j+1,j}$ are displayed as white circles (in the boxes).

As we can see, for generating the system of Eqs. (3.8), i.e., (3.10), we use only entries $s_{k,j}$ for $k \leq n$. Thus, we need the following matrix of the type $(n + 1) \times (2n + 3)$,

$$\mathbf{S} = \begin{bmatrix} s_{0,1} & s_{0,1} & \dots & s_{0,n} & \dots & s_{0,2n+1} & s_{0,2n+2} \\ & s_{1,1} & \dots & s_{1,n} & \dots & s_{1,2n+1} & s_{1,2n+2} \\ & & \ddots & & & & \\ & & & s_{n,n} & \dots & s_{n,2n+1} & s_{n,2n+2} \end{bmatrix}. \tag{5.7}$$

5.3. Method for calculating the nodes

Consider again the system of $n + 1$ nonlinear equations (3.8). Taking only the last n equations of (3.8), we get

$$\mathbf{C}\boldsymbol{\sigma} = \mathbf{d}, \tag{5.8}$$

where $\mathbf{C} = \mathbf{C}(r) = [a_{k,j}]_{k=1,j=1}^{n,n}$ and $\mathbf{d} = \mathbf{d}(r) = [b_1 \cdots b_n]^T$. According to (3.9), the elements $a_{k,j}$ and b_k are expressed in terms of the elements of the matrix \mathbf{S} given by (5.7). The determinant of the matrix \mathbf{C} has the form

$$\det \mathbf{C} = (-1)^{n(n-1)/2} \begin{vmatrix} s_{1,2n} - r s_{1,2n-1} & s_{1,2n-2} - r s_{1,2n-3} & \dots & s_{1,2} - r s_{1,1} \\ s_{2,2n} - r s_{2,2n-1} & s_{2,2n-2} - r s_{2,2n-3} & \dots & s_{2,2} - r s_{2,1} \\ \vdots & \vdots & \ddots & \vdots \\ s_{n,2n} - r s_{n,2n-1} & s_{n,2n-2} - r s_{n,2n-3} & \dots & s_{n,2} - r s_{n,1} \end{vmatrix},$$

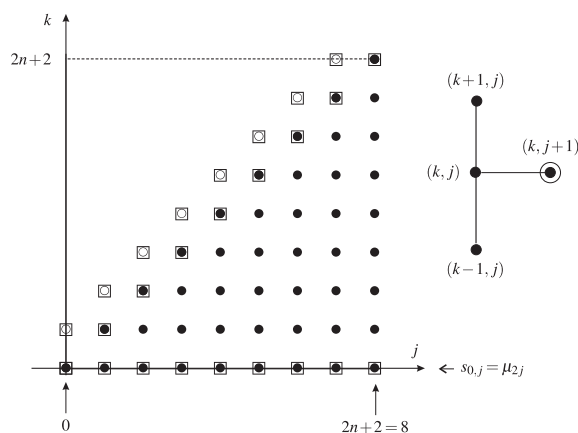


Fig. 2. The scheme for calculating the inner products $s_{k,j} = (\pi_{2k}, z^{2j})$ for $n = 3$.

where its elements $s_{k,2n-2j+2} - rs_{k,2n-2j+1}$ are equal to zero for each $k, j \in \{1, 2, \dots, n\}$ such that $k + 2j > 2n + 2$, because of $s_{k,j} = 0$ for $k > j$. It is easy to see that such zero-elements are the last $n - 2$ elements in the last n th column, the last $n - 4$ elements in the $(n - 1)$ st column, etc. Otherwise, $\det \mathbf{C} = \Delta_n(r)$ is a polynomial of degree n . Since the vector \mathbf{d} on the right side in (5.8) is given by

$$\mathbf{d} = \begin{bmatrix} s_{1,2n+2} - rs_{1,2n+1} \\ s_{2,2n+2} - rs_{2,2n+1} \\ \vdots \\ s_{n,2n+2} - rs_{n,2n+1} \end{bmatrix},$$

the corresponding determinants in Cramer's rule $\Delta_n^{(j)}(r), j = 1, \dots, n$, are also polynomials of degree n . Thus, for a given r , such that $\Delta_n(r) \neq 0$, the unique solution of (5.8) is given by

$$\hat{\sigma}_j = \hat{\sigma}_j(r) = \frac{\Delta_n^{(j)}(r)}{\Delta_n(r)}, \quad j = 1, \dots, n.$$

Using MATHEMATICA package it can be obtained in a symbolic form as rational functions in r . Substituting $\hat{\sigma}_j(r), j = 1, \dots, n$, in

$$f_0 = \sum_{j=0}^n (-1)^j \hat{\sigma}_j \{s_{0,2n-2j+2} - rs_{0,2n-2j+1}\} = 0,$$

we obtain the following algebraic equation of degree $n + 1$,

$$\Phi_{n+1}(r) \equiv \sum_{j=0}^n (-1)^j \{s_{0,2n-2j+2} - rs_{0,2n-2j+1}\} \Delta_n^{(j)}(r) = 0,$$

where $\Delta_n^{(0)}(r) \equiv \Delta_n(r)$.

Numerical experiments show that the (monic) polynomial $\Phi_{n+1}(r)$ has $n + 1$ different real zeros located in $(0, 1)$.

In the case $n = 1$ we have seen in Section 2 that two different solutions exist for r and they give two quadratures of the same precision eleven.

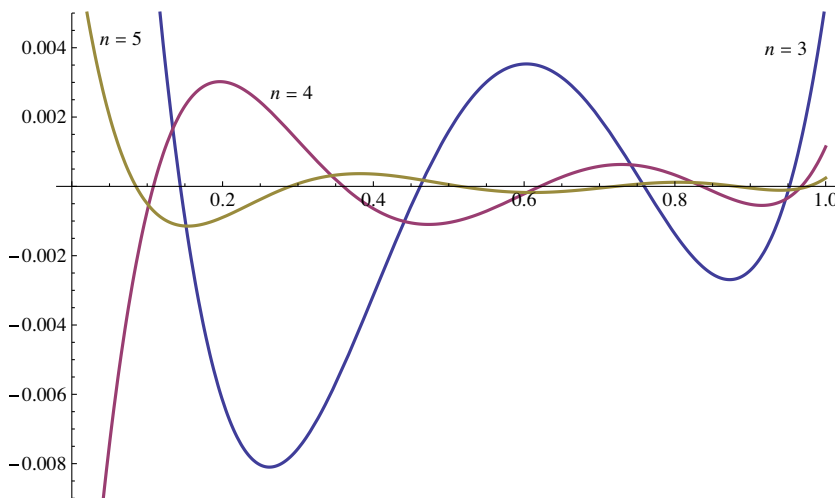


Fig. 3. Graphs of $\Phi_{n+1}(r), n = 3, 4, 5$, for $w(z) = 1$.

Table 1

Different solutions of $\Phi_{n+1}(r) = 0$ for $n = 2(1)5$ and $w(z) = 1$.

$n = 2$	$n = 3$	$n = 4$	$n = 5$
0.2044987378293505	0.1439216162367618	0.1081897446669971	0.08510161904718037
0.6167356745407912	0.4619273121368076	0.3598672165580655	0.2897653961037322
0.9208470355936592	0.7593055545829755	0.6211046569905429	0.5148988061113188
	0.9519663824480733	0.8360221823612692	0.7211387868476094
		0.9678238003414767	0.8814830739148880
			0.9769659264002607

The graphs of monic polynomials $\Phi_{n+1}(r)$ on $[0,1]$ for $n = 3, 4, 5$ are displayed in Fig. 3.

In Table 1 we present the corresponding numerical values of the zeros of the polynomials $\Phi_{n+1}(r)$ for $n = 2, 3, 4, 5$.

Similar results can be obtained for other weight functions. The cases with the Chebyshev weight function of the first kind $w(z) = (1 - z^2)^{-1/2}$ and the weight function $w(z) = |z|$ are displayed in Figs. 4 and 5, respectively.

The previous method for finding all solutions for \hat{r}_0 works for some reasonable values of n , e.g. $n \leq 50$. For example, for the Chebyshev weight of the first kind and $n = 20$ we obtain 21 solutions for \hat{r}_0 (given to 30 decimal digits):

```
Out[7]= {0.0150485864753572668744527330521,
0.0557532669882252298665743917182, 0.108187398514305913909380337895,
0.167824679058336195589651992843, 0.232108607836412242442640888453,
0.299251018982542294927095298222, 0.367850351163304408081805272186,
0.436729680690725174227220507435, 0.504858032671624316488143131170,
0.571309464731063706091020367257, 0.635241286837459266216610097877,
0.695882703617760968552811397194, 0.752529402435763252982400288756,
0.804541615955365152530647156549, 0.851344214028606784515849170776,
0.892427938359564164814715995310, 0.927351213720803665019429733552,
0.955742161852690778718938880015, 0.977300564840333779616628490742,
0.991799603894642762751920785079, 0.9990872539274360180930441192512}
```

In the case $w(z) = 1$ and $n = 50$, the following sequence of 51 solutions for \hat{r}_0 is given with 20 decimal digits:

```
Out[11]= {0.0040059914529631981612, 0.015198148896652530762, 0.030153553927246011838,
0.047786461968069102119, 0.067502176479469289610, 0.088904351714626293601,
0.11170137135198798583, 0.13566492220893164063, 0.16060826297875345727,
0.18637354088540880140, 0.21282382747632203945, 0.23983784595920233875,
0.26730633917663595361, 0.29512949030801918880, 0.32321504715488414291,
0.35147693246320767511, 0.37983419928314226360, 0.40821023692094646313,
0.43653216243154454449, 0.46473035175990477588, 0.49273807747585617047,
0.52049122885459292527, 0.54792809622616881572, 0.57498920592171805202,
0.60161719533987947594, 0.62775672001012940416, 0.65335438628580607294,
0.67835870462602991647, 0.70272005943876842745, 0.72639069223892979455,
0.74932469548418699756, 0.77147801492965236242, 0.79280845872161408075,
0.81327571175331934504, 0.83284135404945289269, 0.85146888214350868901,
0.86912373257358626880, 0.88577330675487417323, 0.90138699659712564769,
0.91593621032745143452, 0.92939439805653787567, 0.94173707669323433174,
0.95294185387170192266, 0.96298845061167595446, 0.97185872249560409663,
0.97953667924407666148, 0.98600850280069667575, 0.99126256482379910043,
0.99528944820940483897, 0.99808200354293588811, 0.99963584456149960413}
```

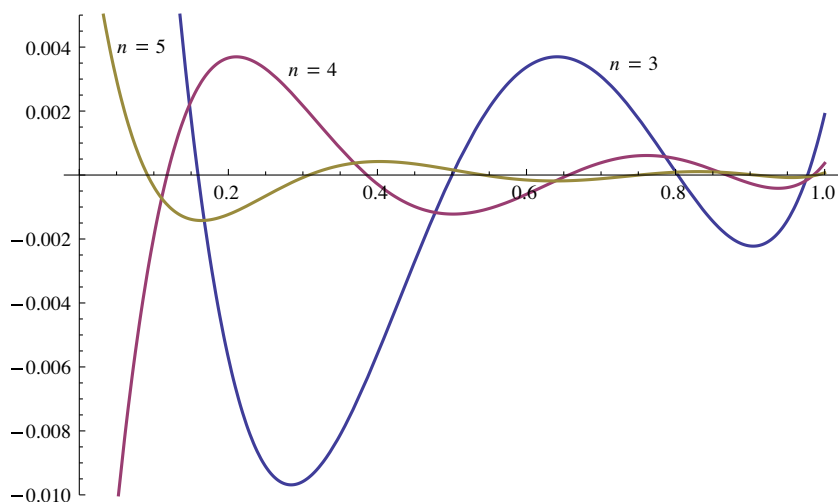


Fig. 4. Graphs of $\Phi_{n+1}(r)$, $n = 3, 4, 5$, for $w(z) = (1 - z^2)^{-1/2}$.

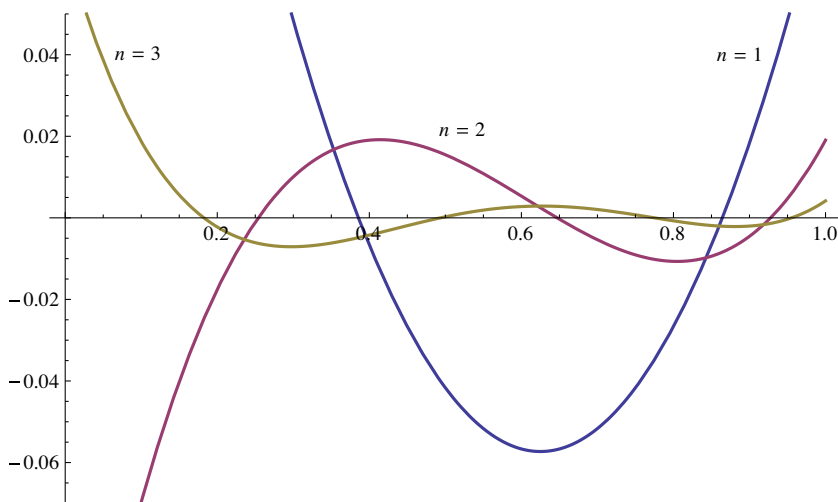


Fig. 5. Graphs of $\Phi_{n+1}(r)$, $n = 1, 2, 3$, for $w(z) = |z|$.

From the theoretical point of view it would be nice to prove (or disprove) the following conjecture:

Conjecture 5.5. For a given weight function $w(z)$ and each $n \in \mathbb{N}$, all zeros of $\Phi_{n+1}(r)$ are real and distinct and are located in $(0, 1)$. Zeros of $\Phi_{n+1}(r)$ and $\Phi_{n+2}(r)$ interlace.

5.4. Numerical results

For quadratures $Q_{4n+3}^M(w; f)$ of degree of precision $d = 6n + 5$, in our MATHEMATICA procedure we obtain complete parameters in the form

$$\{x_0, x_k, A, B, CK, DK\} = \{x_0, \{x_1, \dots, x_n\}, A, B, \{C_1, \dots, C_n\}, \{D_1, \dots, D_n\}\}.$$

For example, let $w(z) = 1$ and $n = 2$. Then we have three quadratures $Q_{11}^M(1; f)$ of degree of precision $d = 17$ (see first column in Table 1). For these values of r_0 , i.e.,

$$\{0.2044987378293505, 0.6167356745407912, 0.9208470355936592\}$$

we get the following parameters:

$$\{0.4522153666444237, \{0.7754684395027309, 0.9570916645968834\}, \\ 0.4880467095490914, 0.3936044844812900, \{0.2527549012554169, 0.1099114468981711\}, \\ \{-0.0003299665322107021, 0.00003577912278703132\}\},$$

$$\{0.7853252030469869, \{0.4741479794169331, 0.9589531260262328\}, 0.5473979047003460, \\ 0.2416521097533237, \{0.3846699903497127, 0.1051207091442720\}, \\ \{-0.005150757567968953, 8.995970487465769 * 10^{-6}\}\},$$

$$\{0.9596077509032840, \{0.4802111190778518, 0.7885463525798828\}, \\ 0.5616568463150571, 0.1034616930531016, \{0.3835087691311978, 0.2383909330938098\}, \\ \{-0.006285348161458679, 0.00009552972582096275\}\},$$

respectively.

In the case of the Chebyshev weight $w(z) = 1/\sqrt{1-z^2}$ and $n = 2$, for the corresponding values of r_0 ,

$$\{0.2321837439443931, 0.6740206457250330, 0.9609831103305740\},$$

the quadrature parameters are respectively:

$$\{0.4818544842007731, \{0.8124087172755511, 0.9790447658917281\}, 0.5249337433901672, \\ 0.4705720970208580, \{0.4268191199148742, 0.4112815014700504\}, \\ \{-0.0003991705374755941, 0.00005590723150597970\}\},$$

Table 2Relative errors in quadrature sums for $w(z) = 1$ and $w(z) = 1/\sqrt{1-z^2}$.

QF	N	d_{\max}	$w(z) = 1$	$w(z) = (1-z^2)^{-1/2}$
$Q_{11,0}^M(w;f)$	4	17	4.44(-5)	7.72(-7)
$Q_{11,1}^M(w;f)$	4	17	5.31(-6)	4.55(-5)
$Q_{11,2}^M(w;f)$	4	17	7.01(-6)	5.79(-5)
$Q_7^G(w;f)$	4	13	2.48(-4)	3.29(-4)
$Q_8^G(w;f)$	4	15	5.73(-6)	3.06(-5)
$Q_9^G(w;f)$	5	17	2.36(-5)	3.67(-5)

{0.8209876038802492, {0.5034904974569647, 0.9799958361402526}, 0.5871100876953411, 0.4159965436705901, {0.4649125011919272, 0.4017086864355387}, {-0.005389674191177676, 0.00001322584034774685}},

{0.9802974601265546, {0.5090243926812192, 0.8235455809669467}, 0.6010918915307779, 0.3986693558288102, {0.4647957620089833, 0.4131917386857352}, {-0.006515105833929775, 0.0001086303399087344}}.

At the end of this paper we give a numerical example. The formula (3.3) could be interesting for real functions of the form $f(z) = g(z^4)$. According to (3.3), in that case each of quadratures $Q_{4n+3,\nu}^M(w;f)$, $\nu = 0, 1, \dots, n$ (for $n+1$ different values of \tilde{r}_0), becomes a quadrature formula of the following form with $N (= n+2)$ nodes,

$$I(w;f) = \int_{-1}^1 f(x)w(x)dx \approx Af(0) + 2Bf(x_0) + \sum_{k=1}^n W_k f(x_k), \quad (5.9)$$

where $W_k = 2(C_k + D_k)$, $k = 1, \dots, n$.

Let $f(x) = 1/(1+x^8)$. We consider two integrals

$$\int_{-1}^1 \frac{dx}{1+x^8} = \frac{1}{4} \left\{ \sin \frac{\pi}{8} \left[\pi + 2 \tanh^{-1} \left(\sin \frac{\pi}{8} \right) \right] + \cos \frac{\pi}{8} \left[\pi + 2 \tanh^{-1} \left(\cos \frac{\pi}{8} \right) \right] \right\} \approx 1.849303411551076,$$

and

$$\int_{-1}^1 \frac{1}{1+x^8} \frac{dx}{\sqrt{1-x^2}} = \pi_4 F_3 \left(\frac{1}{8}, \frac{3}{8}, \frac{5}{8}, \frac{7}{8}; \frac{1}{4}, \frac{1}{2}, \frac{3}{4}; -1 \right) \approx 2.626270969212133.$$

For their calculation we apply all quadratures $Q_{11,\nu}^M(w;f)$, $\nu = 0, 1, 2$, whose parameters are given before. Here $n = 2$, $N = 4$, and $d_{\max} = 17$. The corresponding relative errors

$$\left| \frac{Q_{11,\nu}^M(w;f) - I(w;f)}{I(w;f)} \right|, \quad \nu = 0, 1, 2,$$

are given in the last two columns in Table 2. Numbers in parentheses indicate decimal exponents.

We compare these results with corresponding ones obtained by m -point Gaussian quadratures (for even functions), with respect to Legendre and Chebyshev weights,

$$I(w;f) \approx Q_m^G(w;f) = \sum_{\nu=1}^m A_\nu f(\tau_\nu) = 2 \sum_{\nu=1}^{\lfloor m/2 \rfloor} A_\nu f(\tau_\nu) + \begin{cases} A_{(m+1)/2} f(0), & \text{if } m \text{ is odd,} \\ 0, & \text{if } m \text{ is even,} \end{cases}$$

where A_ν and τ_ν , $\nu = 1, \dots, m$, are Christoffel numbers and nodes, respectively (see [11, pp. 324–325]), such that $1 > \tau_1 > \dots > \tau_m > -1$. If we take $m = 7, 8, 9$, the number of nodes in the corresponding quadrature are $N = 4, 4, 5$, respectively.

As we can see, the quadratures $Q_{11,\nu}^M(w;f)$, $\nu = 0, 1, 2$, have a higher degree of precision than quadratures $Q_m^G(w;f)$, $m = 7, 8$ (with the same number of nodes), as well as that they give a better accuracy.

References

- [1] M. Acharya, B.P. Acharya, S. Pati, Numerical evaluation of integrals of analytic functions, *Internat. J. Comput. Math.* 87 (12) (2010) 2747–2751.
- [2] H. Bateman, A. Erdélyi, *Higher Transcendental Functions, II*, Mc Graw-Hill Book Company, Inc., New York, 1953.
- [3] G. Birkhoff, D.M. Young, Numerical quadrature of analytic and harmonic functions, *J. Math. Phys.* 29 (1950) 217–221.
- [4] T.S. Chihara, *An Introduction to Orthogonal Polynomials*, Gordon and Breach, New York, 1978.
- [5] P.J. Davis, P. Rabinowitz, *Methods of Numerical Integration*, Academic Press, New York, 1975.
- [6] K.V. Laščenov, On a class of orthogonal polynomials, *Učen. Zap. Leningrad. Gos. Ped. Inst.* 89 (1953) 167–189 (Russian).

- [7] F. Lether, On Birkhoff–Young quadrature of analytic functions, *J. Comput. Appl. Math.* 2 (1976) 81–84.
- [8] J.N. Lyness, L.M. Delves, On numerical contour integration round a closed contour, *Math. Comp.* 21 (1967) 561–577.
- [9] J.N. Lyness, C.B. Moler, Numerical differentiation of analytic functions, *SIAM J. Numer. Anal.* 4 (1967) 202–210.
- [10] J.N. Lyness, Quadrature methods based on complex function values, *Math. Comp.* 23 (1969) 601–619.
- [11] G. Mastroianni, G.V. Milovanović, *Interpolation Processes–Basic Theory and Applications*, Springer-Verlag, Berlin–Heidelberg, 2008.
- [12] M.T. McGregor, On a modified Birkhoff–Young quadrature formula for analytic functions, *Univ. Beograd. Publ. Elektrotehn. Fak. Ser. Mat.* 3 (1992) 13–16.
- [13] G.V. Milovanović, Numerical quadratures and orthogonal polynomials, *Stud. Univ. Babeş-Bolyai Math.* 56 (2011) 449–464.
- [14] G.V. Milovanović, A.S. Cvetković, M. Stanić, A generalized Birkhoff–Young–Chebyshev quadrature formula for analytic functions, *Appl. Math. Comput.* 218 (2011) 944–948.
- [15] G.V. Milovanović, R.Ž. Đorđević, On a generalization of modified Birkhoff–Young quadrature formula, *Univ. Beograd. Publ. Elektrotehn. Fak. Ser. Mat. Fiz.* 735–762 (1982) 130–134.
- [16] D.Đ. Tošić, A modification of the Birkhoff–Young quadrature formula for analytic functions, *Univ. Beograd. Publ. Elektrotehn. Fak. Ser. Mat. Fiz.* 602–633 (1978) 73–77.
- [17] D.Đ. Tošić, D. Sotirovski, J. Draškić-Ostojić, Some quadrature formulas for analytic functions, *Univ. Beograd. Publ. Elektrotehn. Fak. Ser. Mat. Fiz.* 678–715 (1980) 170–172.
- [18] D.M. Young, An error bound for the numerical quadrature of analytic functions, *J. Math. Phys.* 31 (1952) 42–44.